

THz generation and transport

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Overview

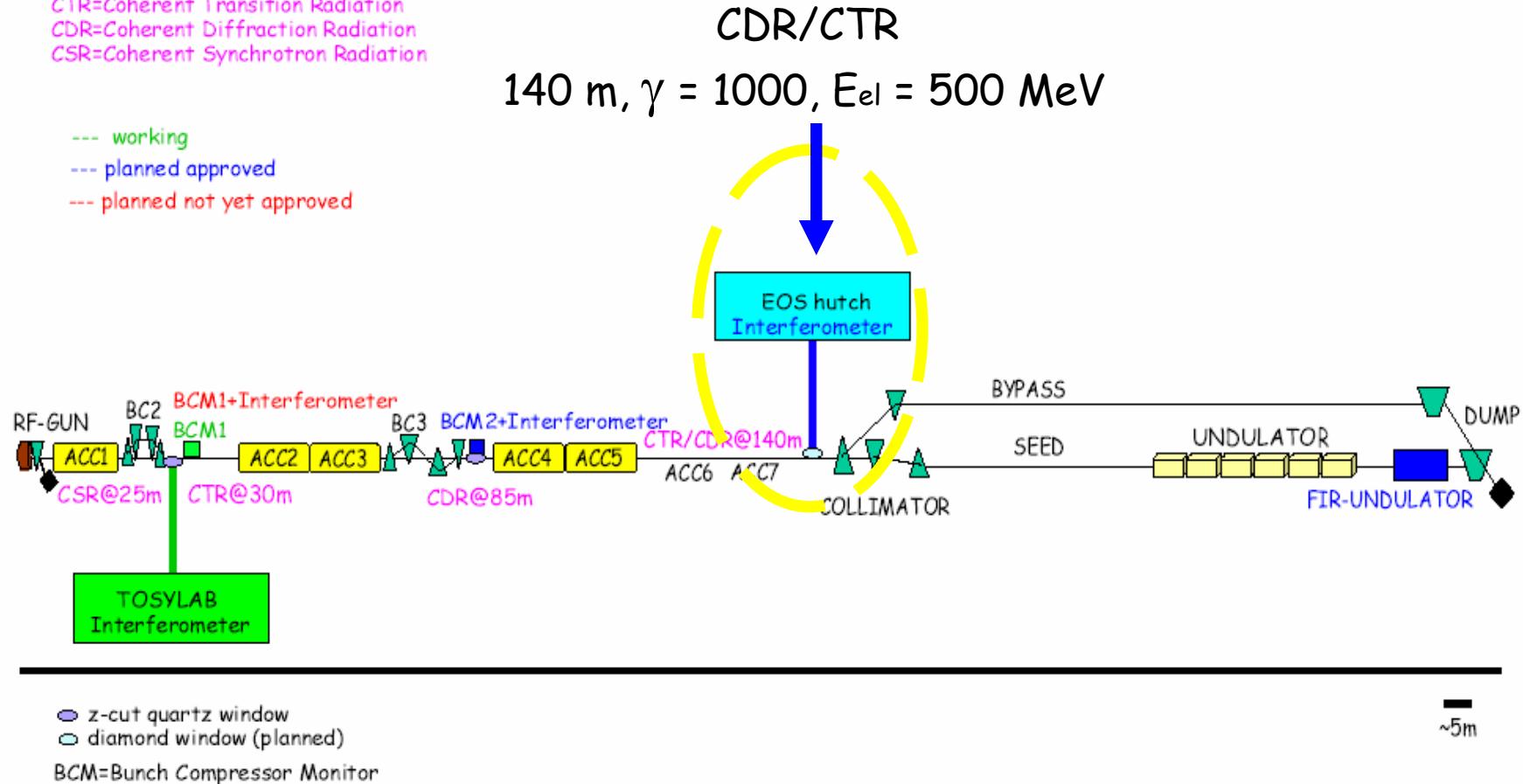
- Motivation
- Design goals
- Simulation tools
 - POP ZEMAX
 - Mathematica code (B. Schmidt, DESY)
- THz beam extraction
- Optical design
- Simulation of the THz radiation transfer line
- Summary and outlook

Motivation

Bunch length measurements with interferometer

CTR=Coherent Transition Radiation
CDR=Coherent Diffraction Radiation
CSR=Coherent Synchrotron Radiation

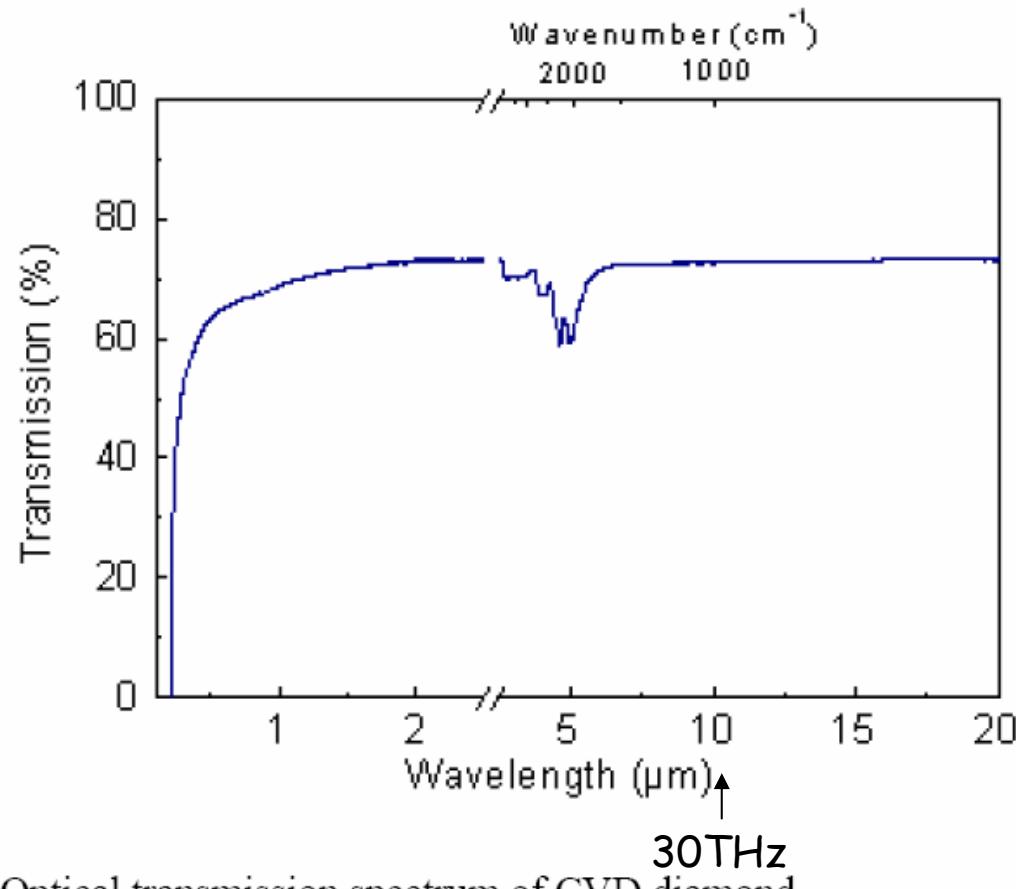
- working
- planned approved
- planned not yet approved



Design goals

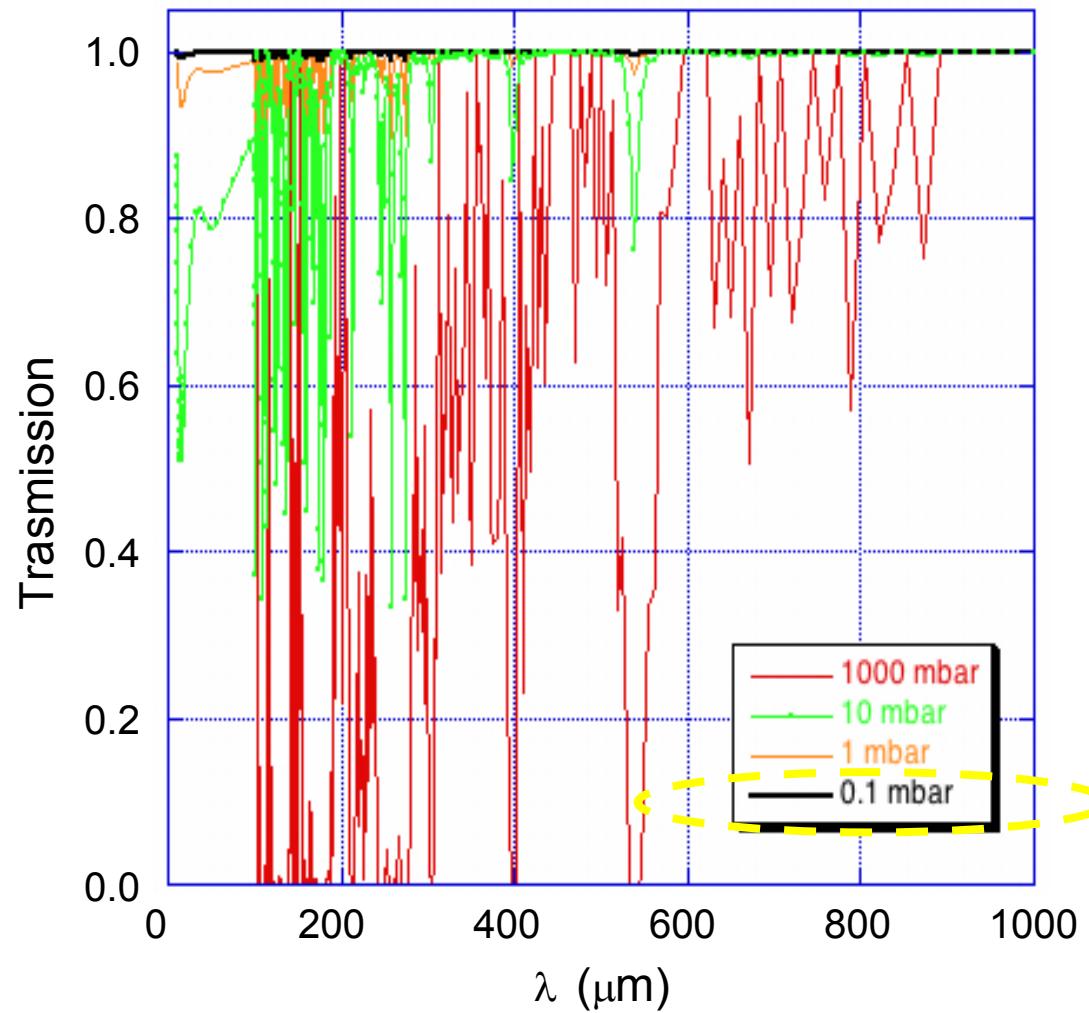
- Flat frequency response
- Low frequency response limit as low as possible BUT
 - Finite dimensions of transfer line tube and mirrors ($\phi \approx 200\text{mm}$)
 - Long transfer line $\sim 20\text{m}$
- High frequency structures expected from μ bunching up to 30THz($10\mu\text{m}$)

Diamond window



Vacuum needed

Transmission through 20m of Humid Air (50% RH)



from B. Schmidt

S. Casalbuoni, DESY

Simulation Tools

2 CODES

ZEMAX

Commercially available

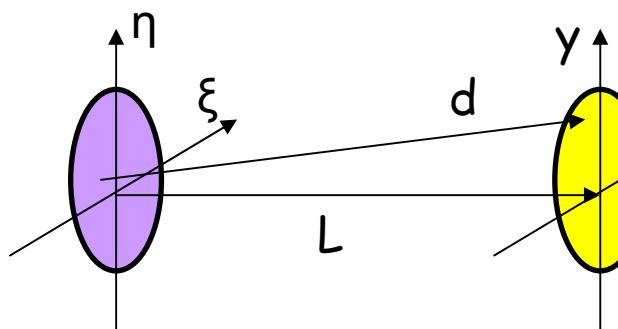
POP (Physical Optic Propagation)

B. Schmidt (DESY)

Mathematica

Huygens-Fresnel principle:

Every point of a wave front may be considered as a centre of a secondary disturbance which gives rise to spherical wavelets, and the wave-front at any later instant may be regarded as the envelope of these wavelets. The secondary wavelets mutually interfere.



Kirchhoff integral

$$\tilde{E}(x, y, L, \omega) = \iint_{\text{Surf}(\xi, \eta)} \tilde{E}(\xi, \eta, 0, \omega) \frac{e^{ikd}}{i\lambda d} d\xi d\eta \approx L$$

finite size screen

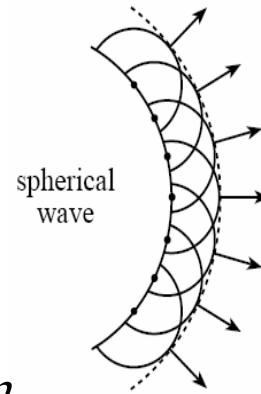
$$d^2 = L^2 + (x - \xi)^2 + (y - \eta)^2; d = L \sqrt{1 + \left(\frac{x - \xi}{L}\right)^2 + \left(\frac{y - \eta}{L}\right)^2}$$

$$|x - \xi|, |y - \eta| \ll L$$

$$d \approx L + \frac{x^2 + y^2}{2L} - \frac{x\xi + y\eta}{L} + \frac{\xi^2 + \eta^2}{2L}$$

**first order
Fraunhofer**

**second order
near field, Fresnel**



S. Casalbuoni, DESY

Input for CTR

Fourier Transform with respect to the longitudinal coordinate $\zeta = z - vt$ of the radial electric field E_r of a uniform bunch charge distribution of radius ρ moving with velocity v in straight line uniform motion (M. Geitz, PhD Thesis).

$$r = \text{radial coordinate} \quad \rho = \text{beam radius} \quad b = \text{pipe radius} \quad k = 2\pi/\lambda = \omega/c$$

$$\tilde{E}_r(k, r) = r \cdot I_1(kr/\gamma) \cdot K_1(k\rho/\gamma) + \rho \cdot I_1(kr/\gamma) \cdot I_1(k\rho/\gamma) \frac{K_0(kb/\gamma)}{I_0(kb/\gamma)}; \quad r < \rho$$

$$\tilde{E}_r(k, r) = \rho \cdot I_1(k\rho/\gamma) \cdot K_1(kr/\gamma) + \rho \cdot I_1(kr/\gamma) \cdot I_1(k\rho/\gamma) \frac{K_0(kb/\gamma)}{I_0(kb/\gamma)}; \quad r > \rho$$

$$\varphi = \text{atan}(y/x)$$

$$\tilde{E}_x = \tilde{E}_r \cdot \cos \varphi \quad \tilde{E}_y = \tilde{E}_r \cdot \sin \varphi$$

Ginzburg-Frank

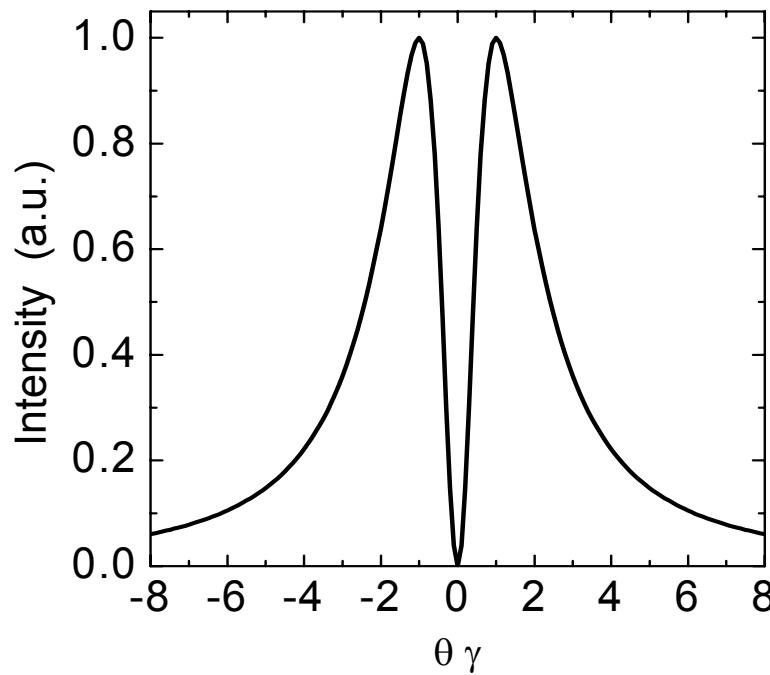
Valid: - if CTR screen radius \gtrsim effective CTR radius

$$a \gtrsim \lambda\gamma$$

-if $L \gg \lambda\gamma^2 \Rightarrow$ far field

(Castellano & Verzilov, Phy.Rev.ST-Accel. Beams ,1998)

$$I(x/L) \propto \frac{\beta^2 \sin^2(x/L)}{(1 - \beta^2 \sin^2(x/L))} \quad \frac{x}{L} = \theta \Rightarrow \text{frequency independent}$$



$\gamma = 100$

radius CTR screen [m]:

 $a = 30\text{mm}$ distance to observation screen [m]: $L = 10\text{m}$

wavelength

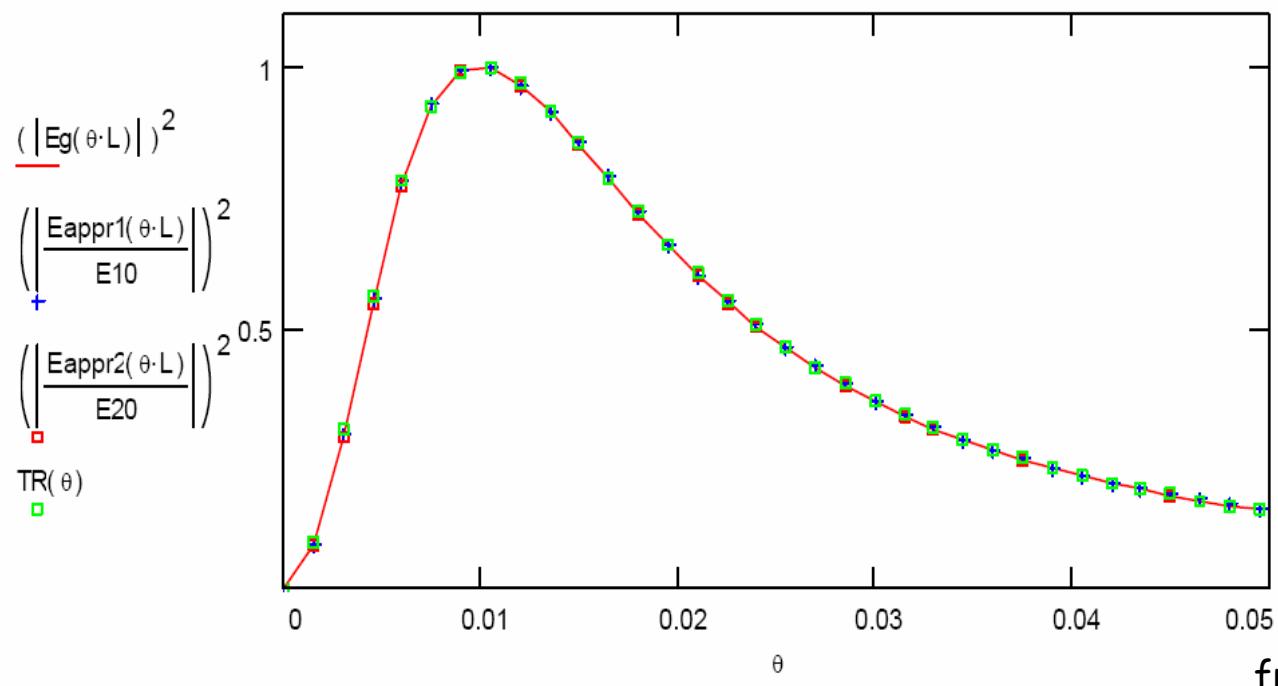
 $\lambda = 300\mu\text{m}$

$$\gamma \frac{\lambda}{a} = 1$$

$$\frac{L}{\gamma^2 \cdot \lambda} = 3.333$$

Both conditions satisfied

green: Ginzburg-Frank | blue: first order, red: exact SQRT resp. second order



from P. Schmüser

S. Casalbuoni, DESY

$\gamma = 100$ radius CTR screen [m]: $a = 30\text{mm}$ distance to observation screen [m]: $L = 20\text{m}$

wavelength

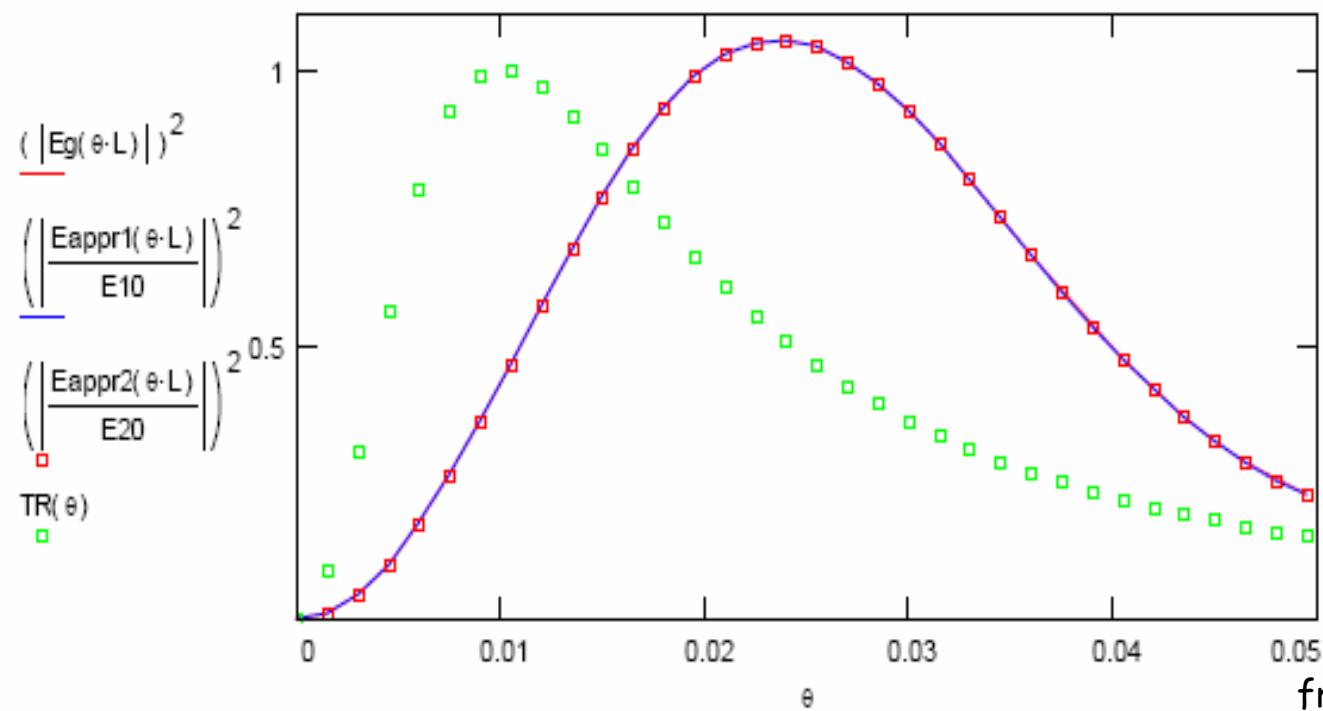
**Finite size
CTR screen**

$$\gamma \frac{\lambda}{a} = 5$$

$$\frac{L}{\frac{a^2}{\gamma \cdot \lambda}} = 1.333$$

 $\lambda = 1.5\text{mm}$

green: Ginzburg-Frank, blue: first order, red: exact SQRT resp. second order



S. Casalbuoni, DESY

$\gamma = 100$ radius CTR screen [m]: $a = 30\text{mm}$ distance to observation screen [m]: $L = 0.5\text{m}$

wavelength

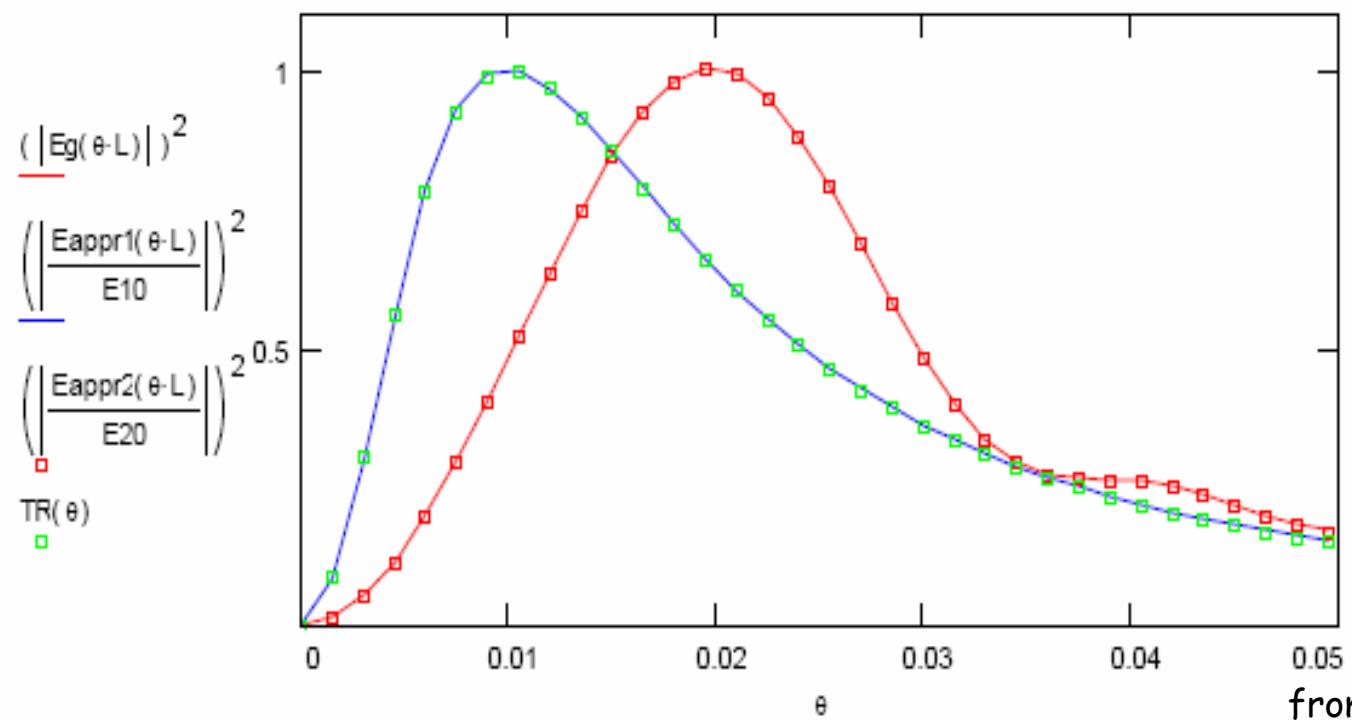
$$\gamma \frac{\lambda}{a} = 1$$

$$\frac{L}{\frac{2}{\gamma} \cdot \lambda} = 0.167$$

$$\lambda = 300\mu\text{m}$$

Near field

green: Ginzburg-Frank, blue: first order, red: exact SQRT resp. second order



from P. Schmüser

S. Casalbuoni, DESY

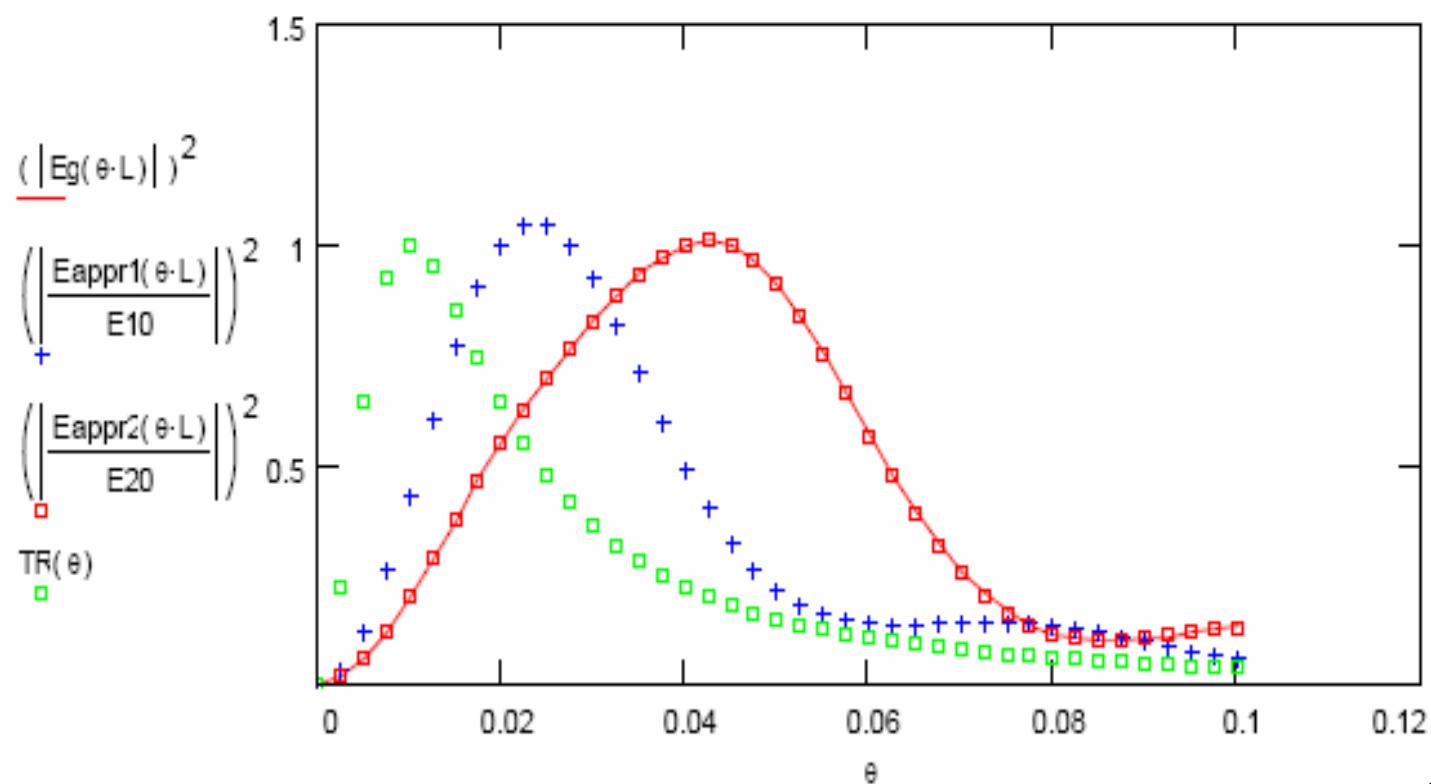
$\gamma = 100$ $L = 0.5\text{m}$ $a = 30\text{mm}$ $\lambda = 1.5\text{mm}$

$$\frac{\lambda \cdot \gamma}{a} = 5$$

$$\frac{L}{\lambda \cdot \gamma^2} = 0.033$$

Both conditions violated

green: Ginzburg-Frank, blue first order red: second order and exact SQRT



from P. Schmüser

S. Casalbuoni, DESY

Simulation Tools

2 CODES

ZEMAX

Commercially available

POP (Physical Optic Propagation)

B. Schmidt (DESY)

Mathematica

Both make use of scalar Fresnel diffraction theory

Valid if objects and apertures $\gg \lambda$

$|x-\xi|, |y-\eta| \ll L$

$$\tilde{E}(x, y, L, \omega) = \frac{1}{i\lambda d} \iint_{Surf(\xi, \eta)} \tilde{E}(\xi, \eta, 0, \omega) e^{ikd} d\xi d\eta$$

Fourier transform of

$$\tilde{E}(x, y, L, \omega) = \frac{1}{i\lambda L} \iint_{Surf(\xi, \eta)} \tilde{E}(\xi, \eta, 0, \omega) e^{ik(\xi^2 + \eta^2)/2L} e^{ik(x\xi + y\eta)/2L} d\xi d\eta$$

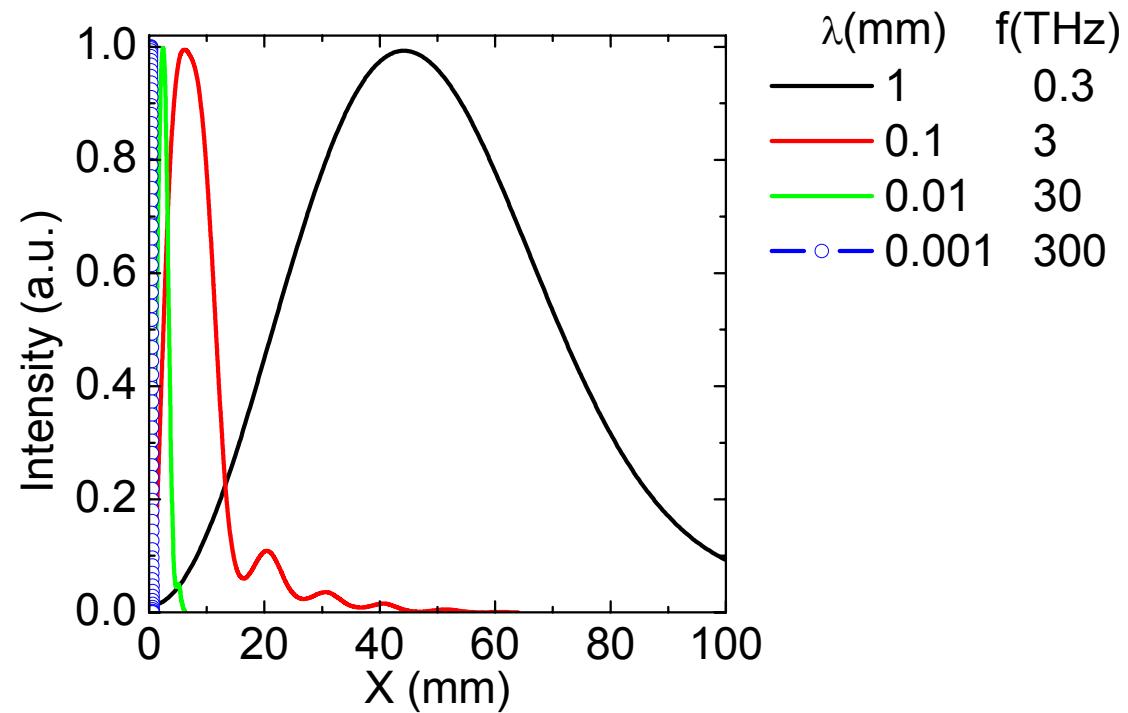
$$k = \frac{2\pi}{\lambda}; \quad \lambda = \frac{2\pi c}{\omega}$$

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Free propagation

if $a \lesssim \gamma\lambda$ and/or $L \lesssim \lambda/\gamma^2 \Rightarrow I(x, L, \omega)$ frequency dependent

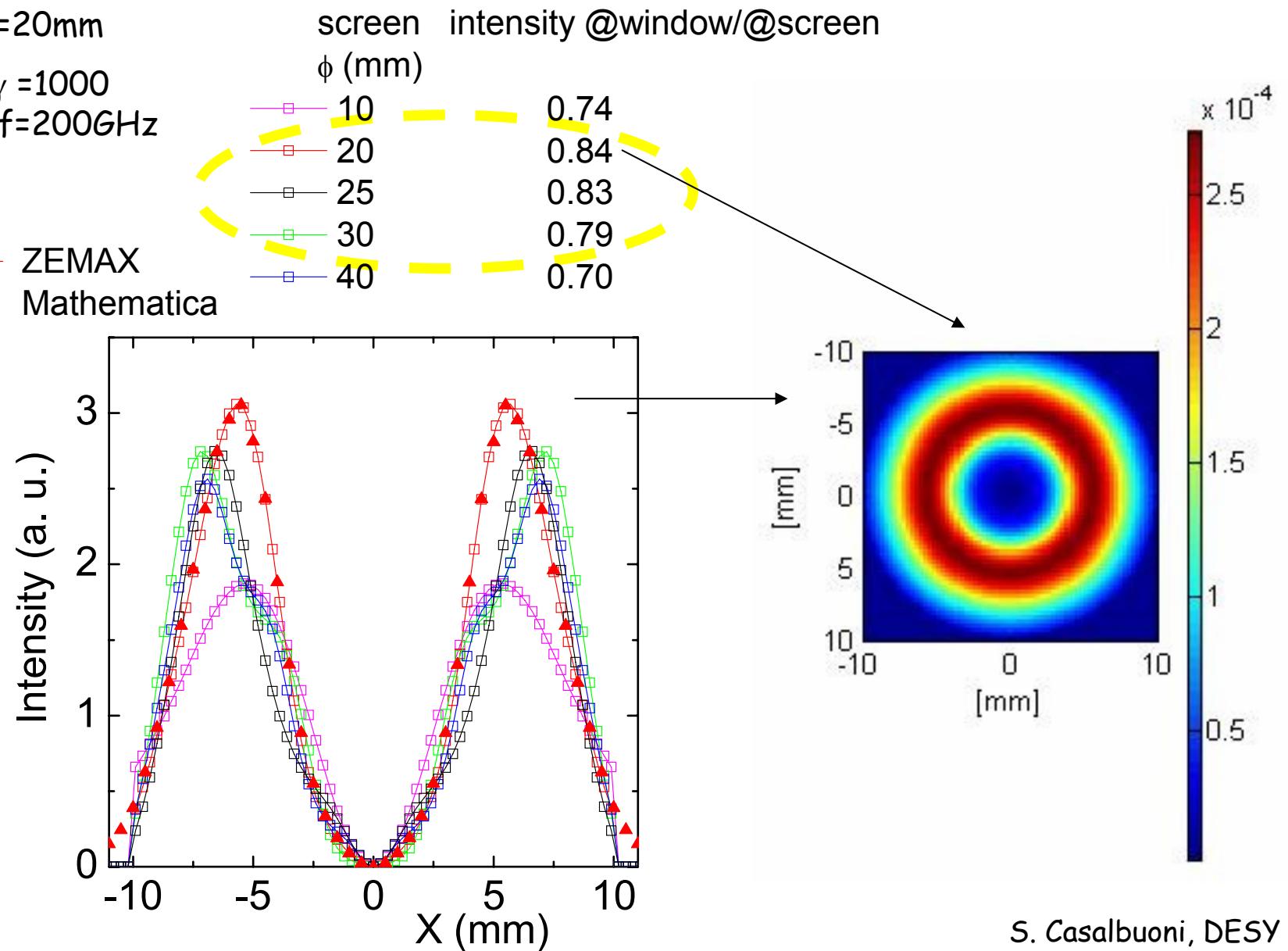
$L=1$ m; $a=10$ mm; $\gamma=1000$



Dimensions of the CTR screen

window $\phi=20\text{mm}$
 $L=40\text{mm}; \gamma=1000$
 $\lambda=1.5\text{mm} \Rightarrow f=200\text{GHz}$

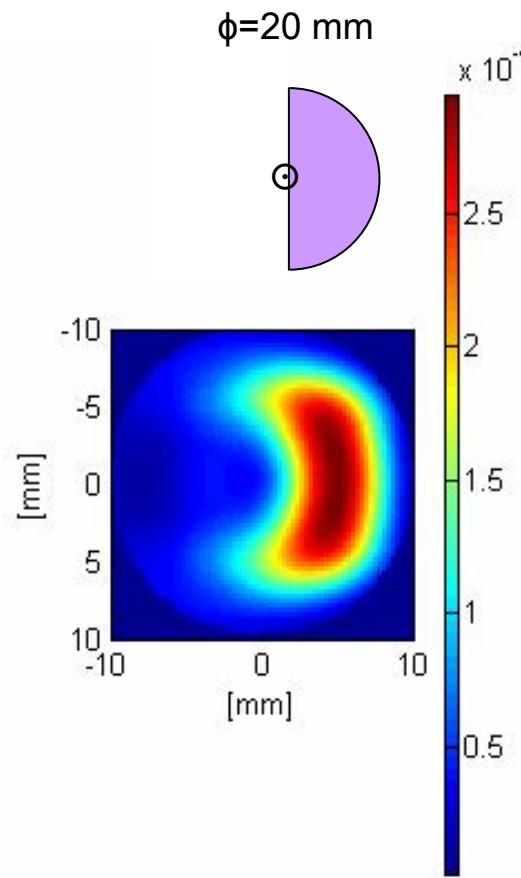
ZEMAX
Mathematica



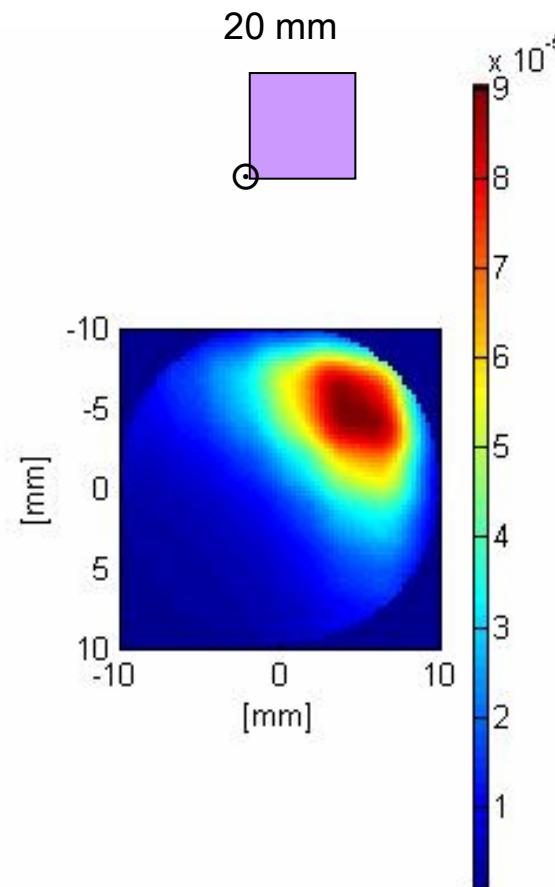
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window $\phi=20\text{mm}$
 $L=40\text{mm}; \gamma =1000$
 $\lambda=1.5\text{mm} \Rightarrow f=200\text{GHz}$

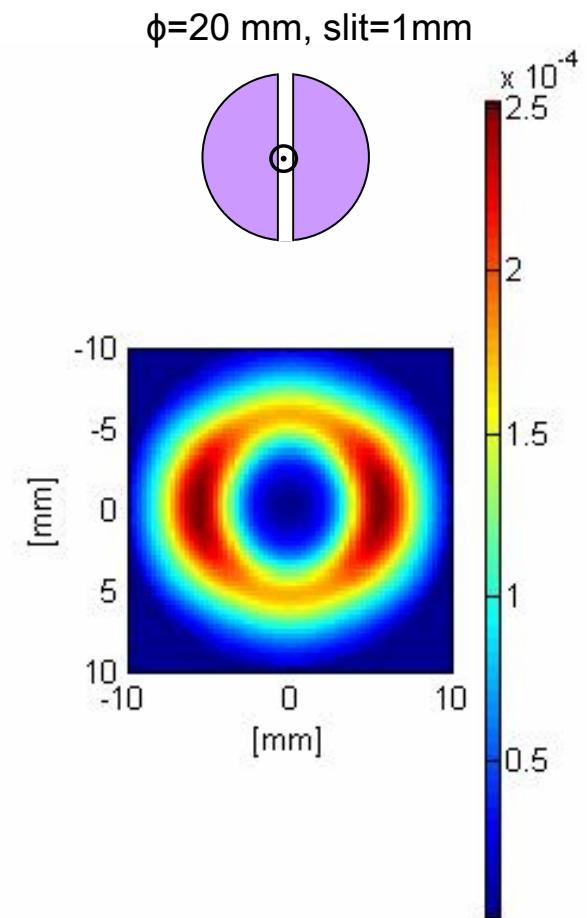
CDR screen



intensity
@window/@screen
0.75

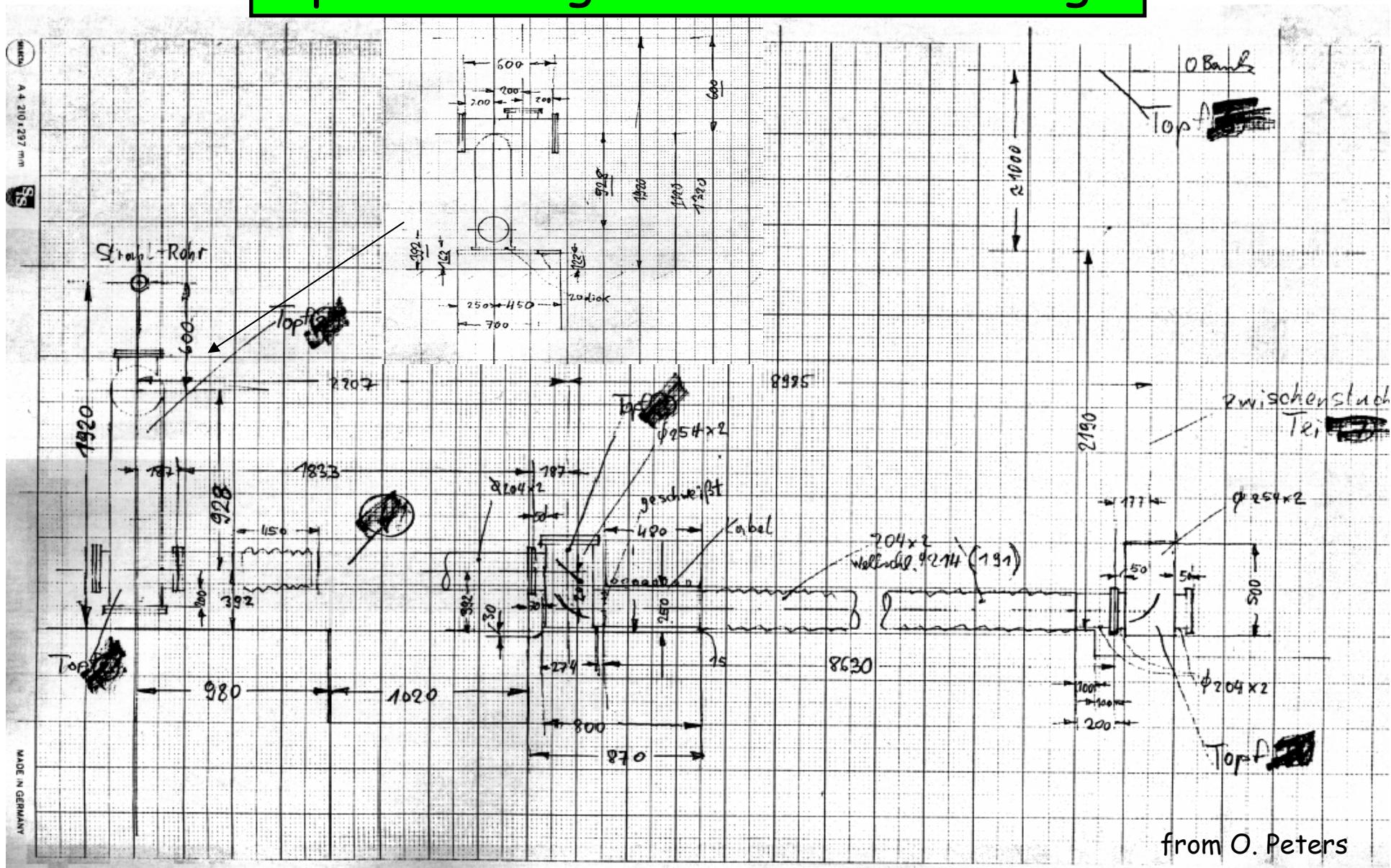


0.55



0.78

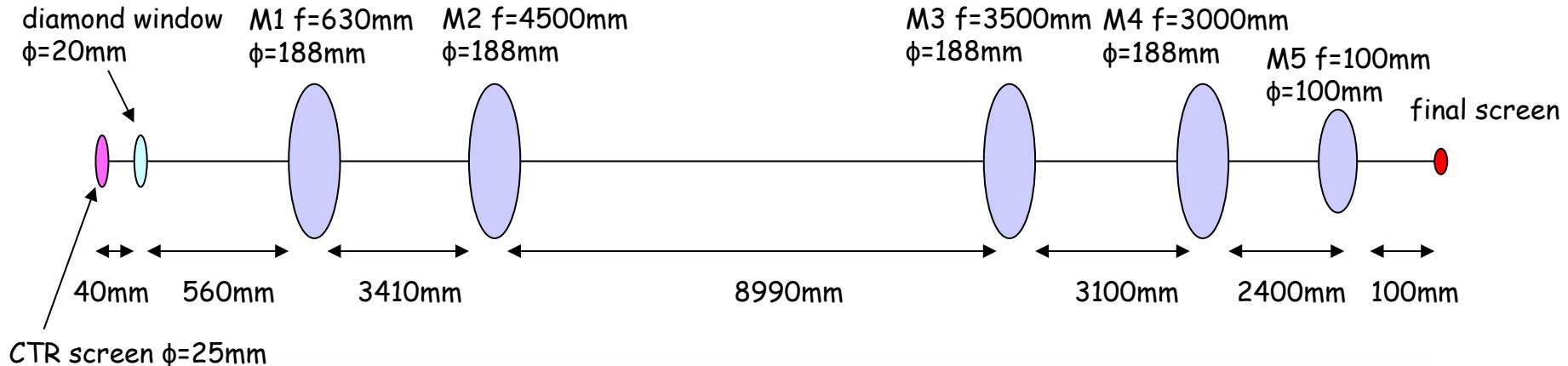
Optical design: technical drawing



from O. Peters

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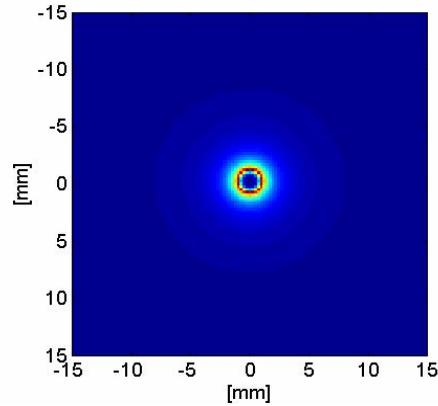
Optical design



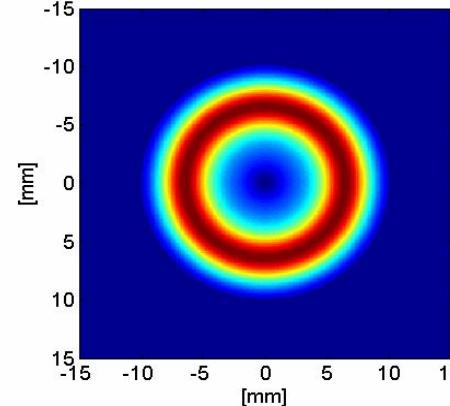
Simulation of the THz radiation transfer line with ideal thin lenses

$\lambda=1.5\text{mm} \Rightarrow f=200\text{GHz}; \gamma = 1000$

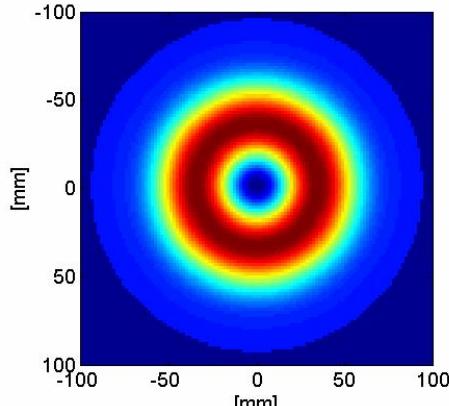
CTR screen $\phi=25\text{mm}$



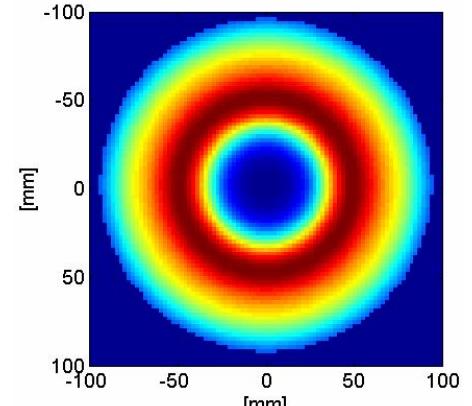
Diamond window $\phi=20\text{mm}$



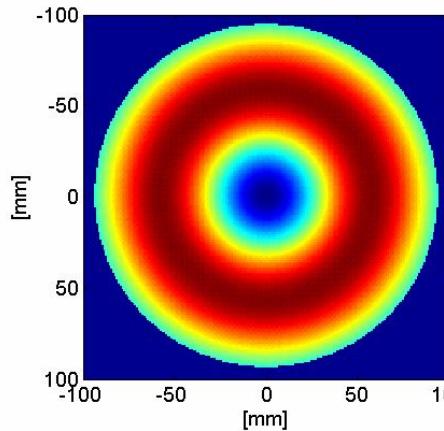
M1



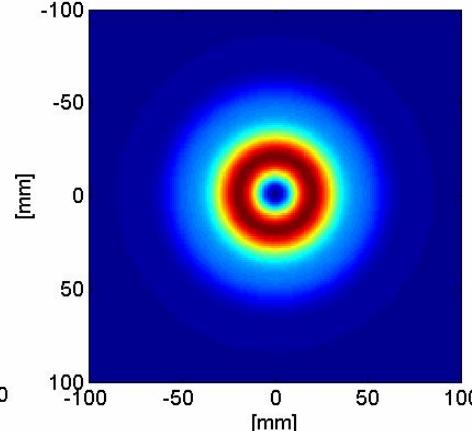
M2



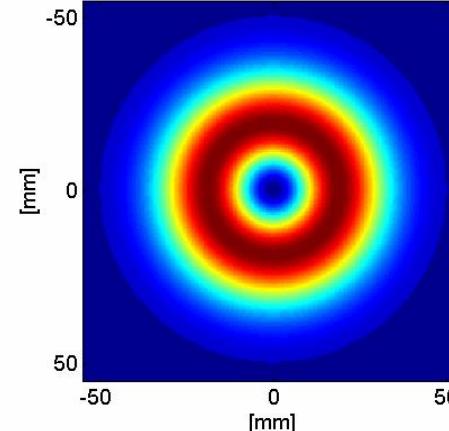
M3



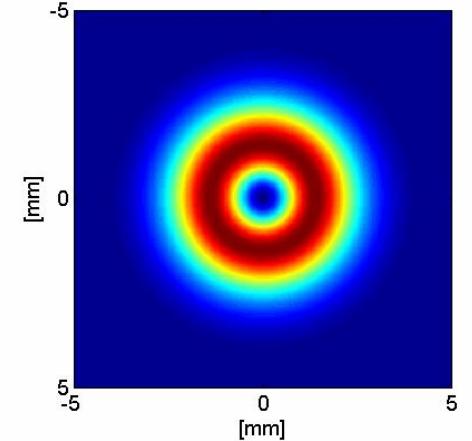
M4



M5



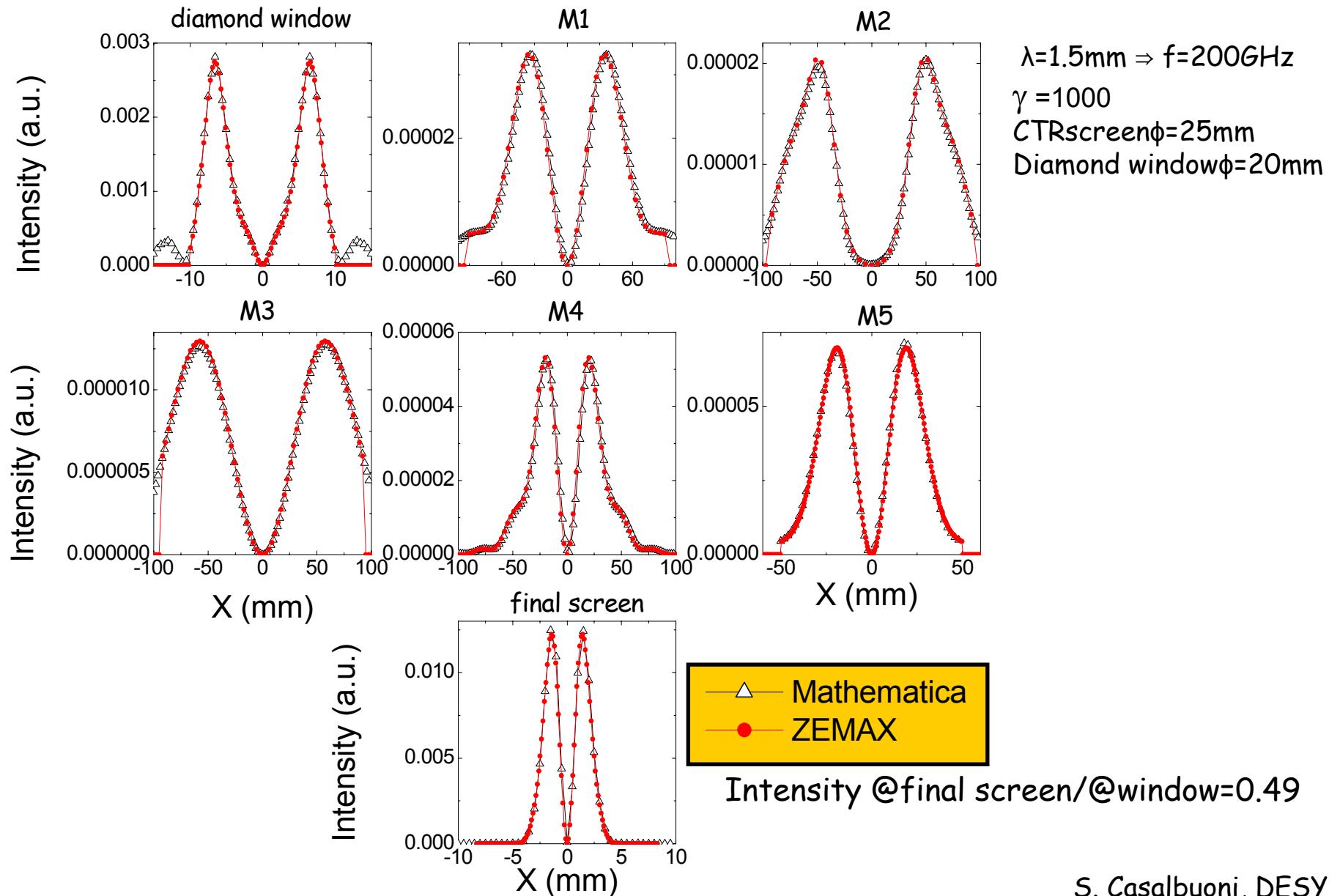
final screen $\phi=8\text{mm}$



Intensity @final screen/@window=0.49

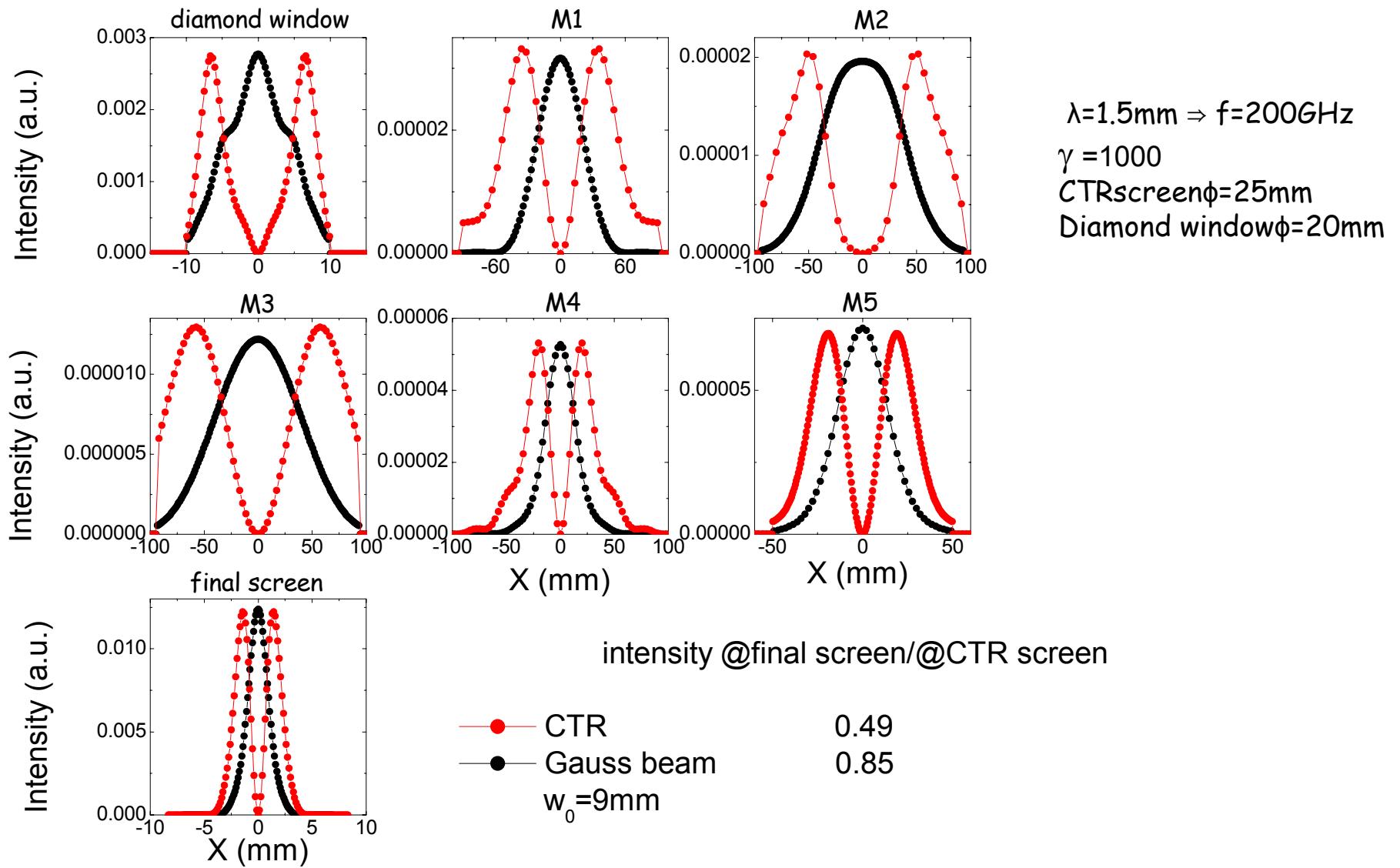
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Simulation of the THz radiation transfer line with ideal thin lenses

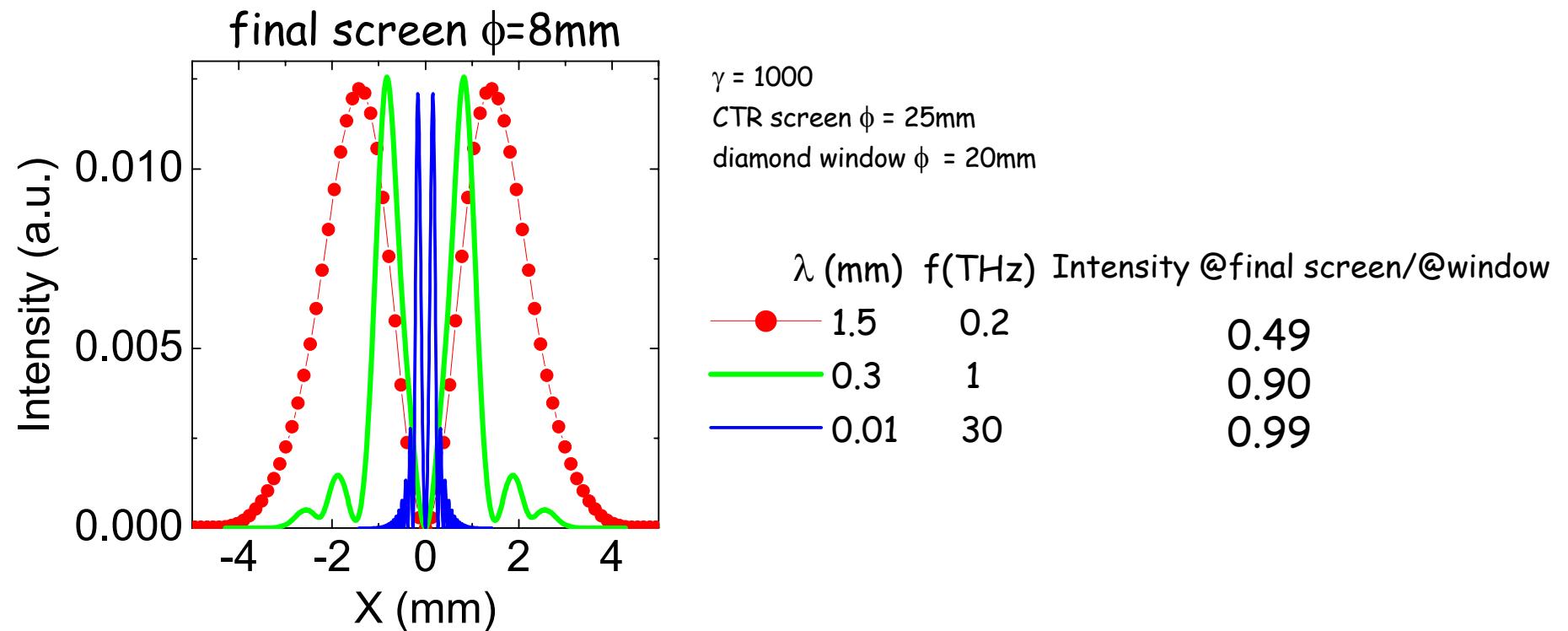


S. Casalbuoni, DESY

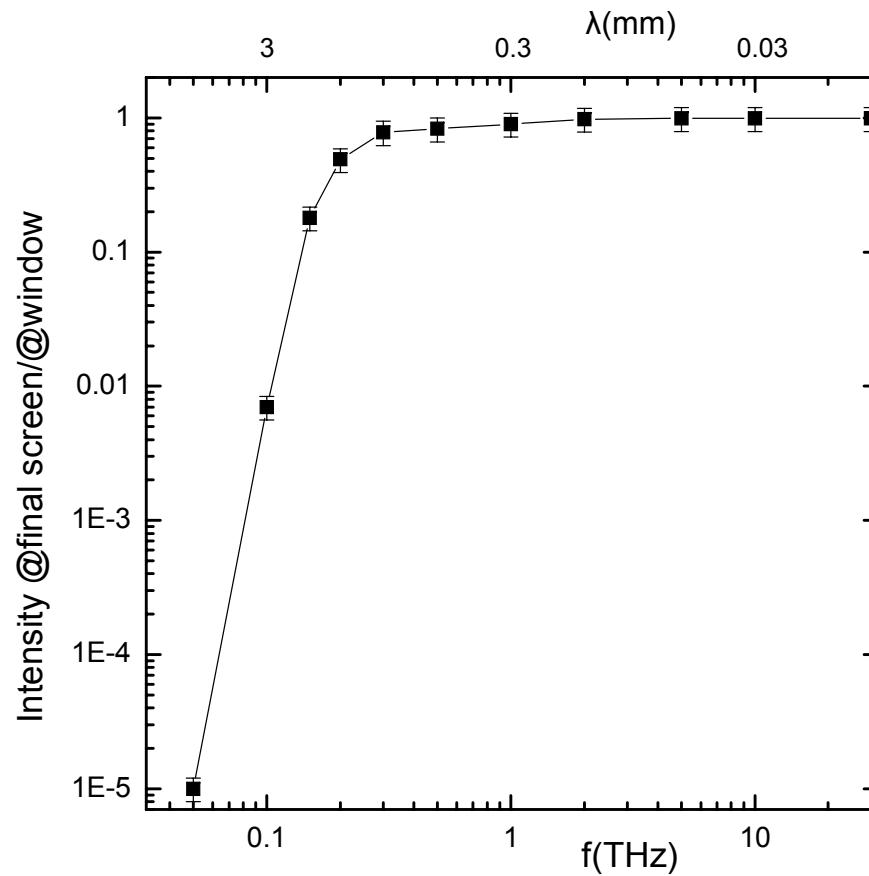
Good also for a Gaussian beam



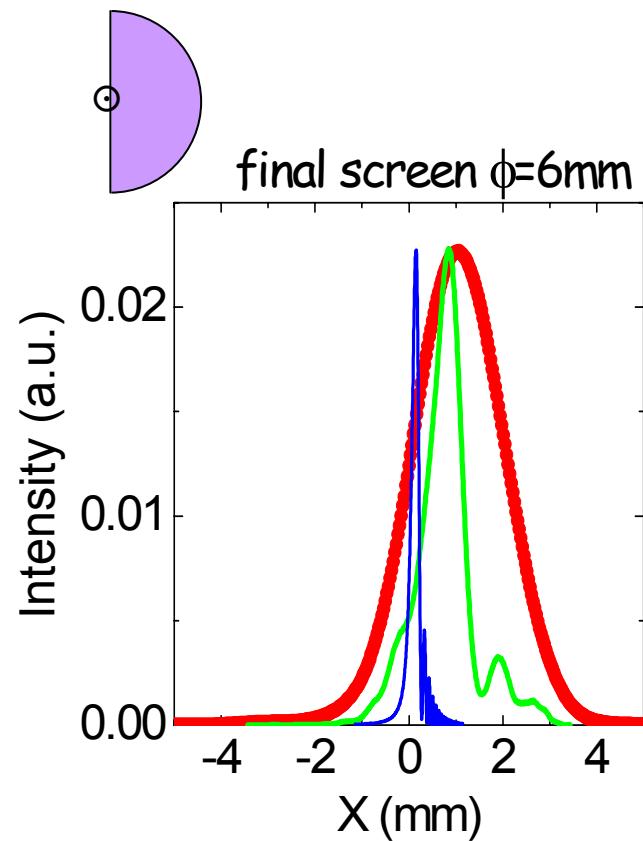
Simulation of the THz radiation transfer line at different frequencies



Transfer function

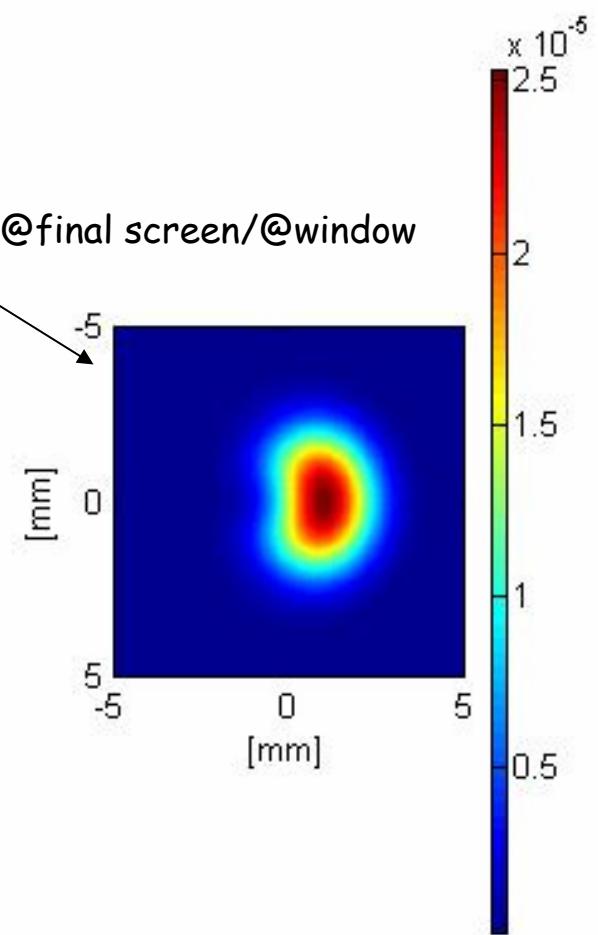


Half screen and horizontal polarization: frequency response

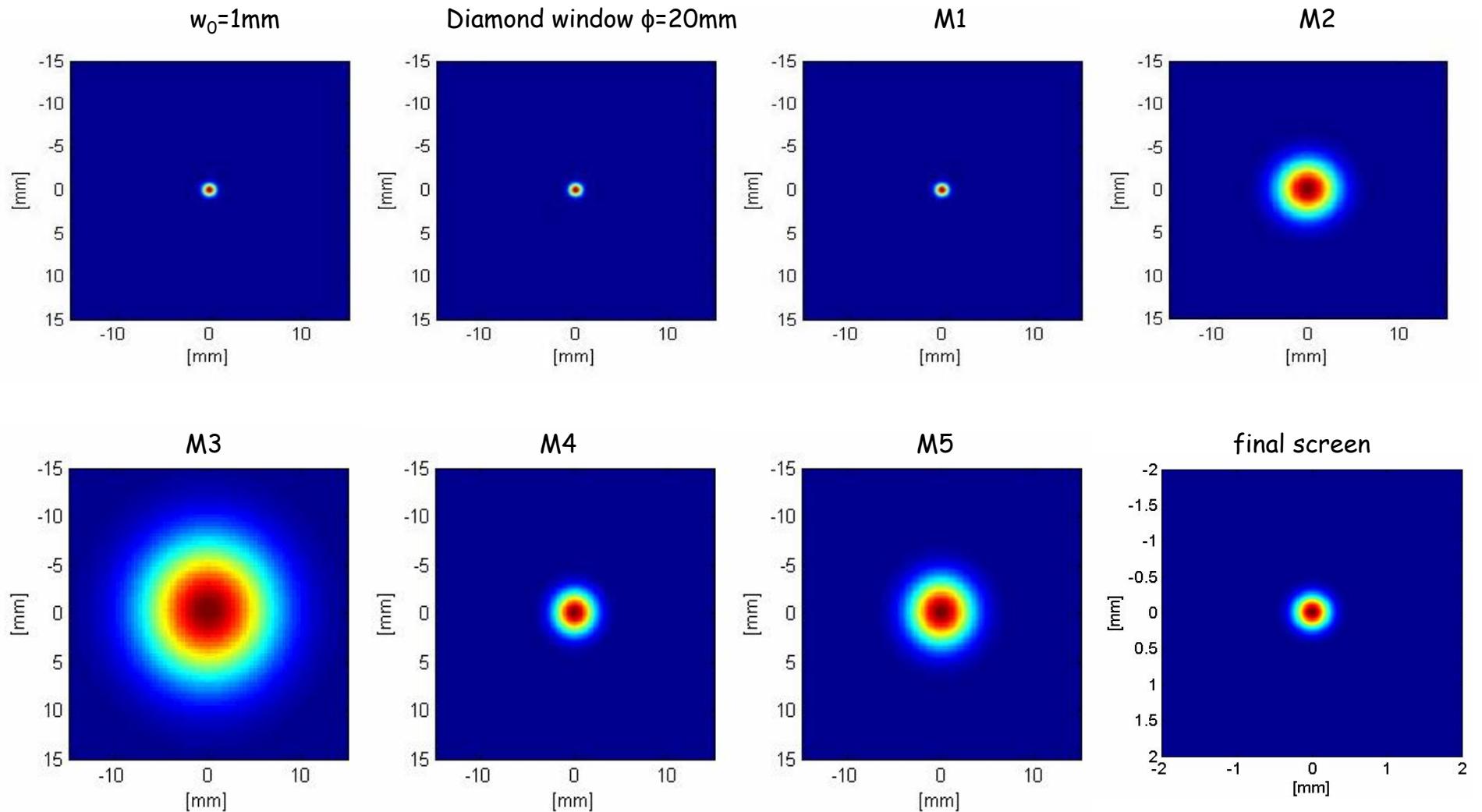


$\gamma = 1000$
CTR screen $\phi = 25\text{mm}$
diamond window $\phi = 20\text{mm}$

λ (mm)	f(THz)	Intensity @final screen/@window
1.5	0.2	0.52
0.3	1	0.91
0.01	30	0.99



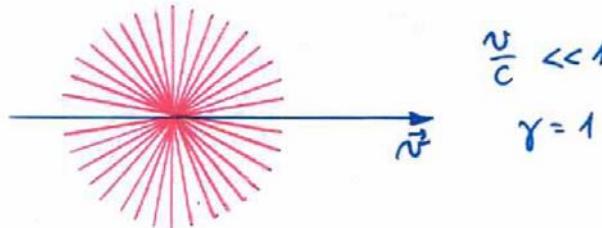
Gaussian beam $\lambda=500\text{nm}$



Summary and outlook

- 2 simulation tools: very good agreement
- Optical design for the THz beam line transfer @140m in TTF2: tested for CTR, CDR and Gaussian beam
- Outlook
 - "ideal" mirror surface
 - tests for stability against beam displacement and mirrors misalignment
 - effect of tilting CTR screen

Electric field of a charge in the laboratory system



$$\vec{E}(\vec{r}) = \frac{q}{4\pi\epsilon_0} \cdot \frac{(1 - \beta^2)}{(1 - \beta^2 \sin^2 \theta)^{3/2}} \cdot \frac{\vec{r}}{r^3}$$

$$E_z \equiv 0; \quad E_r = \frac{q}{4\pi\epsilon_0} \cdot \gamma \cdot \frac{r}{[\gamma^2(z - vt)^2 + r^2]^{3/2}}$$

$$\zeta = z - vt = z - \beta ct$$

$$\tilde{E}_r(k, r) = \int_{-\infty}^{\infty} E_r(k, r) \cdot e^{ik\zeta} d\zeta$$

$$\tilde{E}_r(k, r) = \frac{q}{4\pi\epsilon_0} \cdot \frac{k}{\gamma} \cdot K_1\left(\frac{kr}{\gamma}\right)$$

